

Resonant monomial

(15)

* In the presence of resonance, e.g., $\lambda_i = m_1 \lambda_1 + \dots + m_n \lambda_n$
the monomial term $x_1^{m_1} \dots x_n^{m_n} e_i$ cannot be cancelled
via polynomial change of variable: resonant monomial

* Adapting the procedure of the Poincaré-Dulac Theorem
we expect to be able to reduce a dynamical system to a
form:

$$\dot{X} = LX + P_n(X) + O(|X|^{n+1})$$

where $P_n(X)$ is the sum of resonant monomials up to order n .

* Such a form only depends on the linear part $L = DF(0)$
(considered under its Jordan form) and is called
a normal form.

Normal form theorem

$$\dot{X} = F(X), \quad F \text{ of class } C^k, \quad F(0) = 0, \quad DF(0) = L$$

Choose a complement G_k such that $G_k + \text{ad}_L H_k = H_k$.

There is an analytical change of coordinates such
that in a neighborhood of the origin the system

$\dot{X} = F(X)$ is transformed into:

$$\dot{Y} = LY + Q_2(Y) + \dots + Q_r(Y) + o(|Y|^r), \quad Q_i \in G_i, \\ 2 \leq i \leq r$$

(16)

Normal form for $L = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \Leftrightarrow \begin{cases} \dot{x} = -y + o(|x|, |y|) \\ \dot{y} = x + o(|x|, |y|) \end{cases}$

* $Ad_L: H_2 \rightarrow H_2$:

$$\begin{aligned} Ad_L \begin{pmatrix} x^2 \\ 0 \end{pmatrix} &= \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} x^2 \\ 0 \end{pmatrix} - \begin{pmatrix} 2x & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \\ &= \begin{pmatrix} 0 \\ x^2 \end{pmatrix} - \begin{pmatrix} -2xy \\ 0 \end{pmatrix} = \begin{pmatrix} 2xy \\ x^2 \end{pmatrix} \end{aligned}$$

$\begin{pmatrix} x^2 \\ 0 \end{pmatrix} = x^2 e_x$. basis: $x^2 e_x \quad xy e_x \quad y^2 e_x$
 $x^2 e_y \quad xy e_y \quad y^2 e_y$

basis \rightarrow

Matrix:

$$\begin{array}{c} \text{basis} \downarrow \\ \left[\begin{array}{ccc|ccc} 0 & -1 & 0 & -1 & 0 & 0 \\ 2 & 0 & -2 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 & 0 & -1 \\ \hline 1 & 0 & 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 2 & 0 & -2 \\ 1 & 0 & 1 & 0 & 1 & 0 \end{array} \right] \end{array}$$

\rightarrow invertible: all quadratic terms can be removed

* Evaluating the matrix of $Ad_L: H_3 \rightarrow H_3$ reveals that the resonant terms $(x^2+y^2)(x e_x + y e_y)$ and $(x^2+y^2)(-y e_x + x e_y)$ cannot be removed.

Thus the system can be transformed to:

$$\begin{cases} \dot{x} = -y + (x^2+y^2)(ax-by) + o(|x|, |y|^3) \\ \dot{y} = x + (x^2+y^2)(ay+bx) + o(|x|, |y|^3) \end{cases}$$

Normal form for Hopf bifurcation

* Normal form for $L = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$ in polar coordinate

$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta \end{cases} \Rightarrow \begin{cases} \dot{r} = ar^3 + o(|r|^3) \\ \dot{\theta} = 1 + br^2 + o(|r|^3) \end{cases} \quad \begin{array}{l} r=0 \text{ stable} \\ \text{if } a > 0. \end{array}$$

Hopf bifurcation

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = F(x, y; \mu) = L_{\mu} \begin{pmatrix} x \\ y \end{pmatrix} + G(x, y; \mu), \quad \mu = 0$$

$$F(0, 0; 0) = 0, \quad DF(0, 0; \mu) = L_{\mu}, \quad L_0 = \begin{pmatrix} \mu & -w \\ w & \mu \end{pmatrix}$$

$$G(0, 0; 0) = 0, \quad DG(0, 0; 0) = 0$$

From normal form analysis; the above system can be reduced to $G(x, y; 0) = (x^2 + y^2) \begin{pmatrix} ax - by \\ ay + bx \end{pmatrix}$

Canonical model (normal form for Hopf bifurcation):

$$\dot{x} = \mu x - wy + (x^2 + y^2)(ax - by) + o(|x, y|^3)$$

$$\dot{y} = wx + \mu y + (x^2 + y^2)(ay + bx) + o(|x, y|^3)$$

In complex notations: $z = x + iy$

$$\dot{z} = (\mu + iw)z + (a - ib)|z|^2 z + o(|z|^3)$$

⌊ Truncated normal form is a topological normal form: no need for higher terms;

Codimension of bifurcation

- * The notion of codimension is related to the topological notion of transversality (not discussed).
- * Informally, the codimension of a bifurcation refers to the number of conditions (equalities on the bifurcation parameters) that need to be met at the bifurcation points:
- * Local bifurcation occurs at nonhyperbolic equilibria for which the linearization L of a system $\dot{x} = F(x)$ has eigenvalue with zero real parts. The codimension is the number of eigenvalues with zero real part (conjugate eigenvalues count for one)

* Codimension 1 bifurcation corresponds to

$$L = \left[\begin{array}{c|c} 0 & 0 \\ \hline 0 & A \end{array} \right]$$

saddle - node
 =
 real eigenvalue
 crosses the imaginary axis

or

$$L = \left[\begin{array}{c|c} 0 - \omega & 0 \\ \omega & 0 \\ \hline 0 & A \end{array} \right]$$

Hopf bifurcation
 =
 complex conjugate
 eigenvalue cross the
 imaginary axis

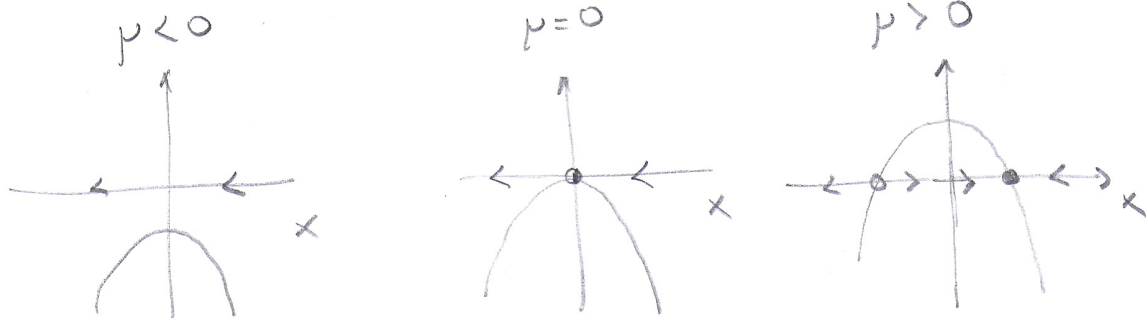
* Codimension 2:

$$L = \left[\begin{array}{c|c} \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} & 0 \\ \hline 0 & A \end{array} \right], \left[\begin{array}{c|c} 0 - \omega & 0 & 0 \\ \omega & 0 & 0 \\ 0 & 0 & 0 \\ \hline 0 & A \end{array} \right], \left[\begin{array}{c|c} 0 - \omega_1 & 0 & 0 & 0 \\ \omega_1 & 0 & 0 & 0 \\ 0 & 0 & 0 & \omega_2 \\ 0 & 0 & -\omega_2 & 0 \\ \hline 0 & A \end{array} \right]$$

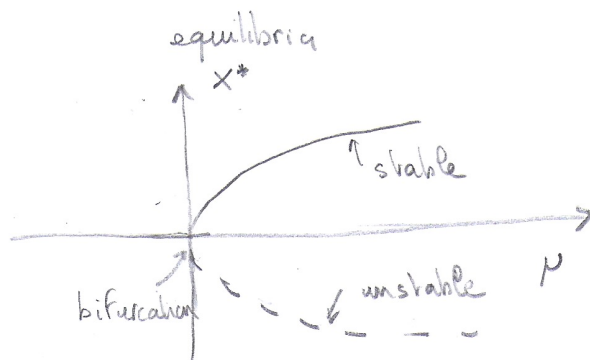
Codimension-1 local bifurcation

* saddle-node (fold) bifurcation

↳ "proper" saddle node: normal form $\dot{x} = \mu - x^2$



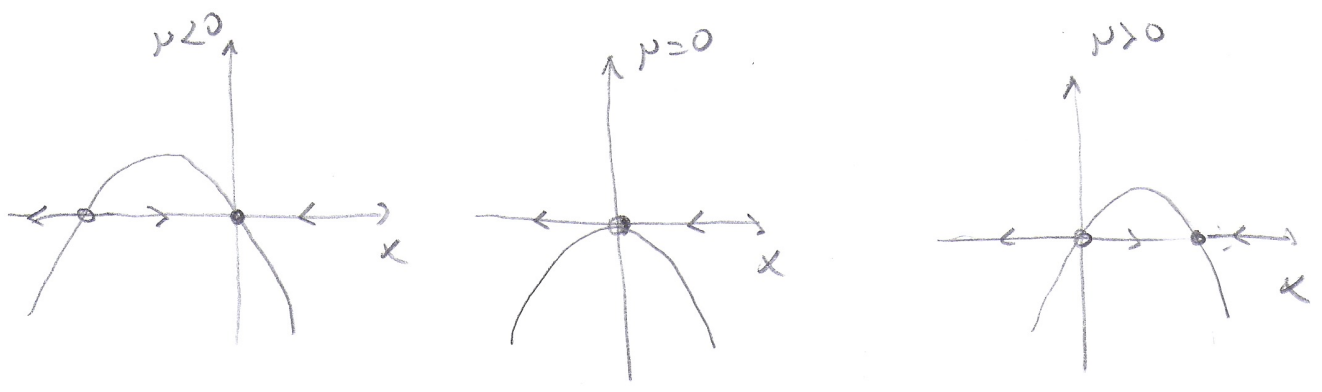
bifurcation
diagram



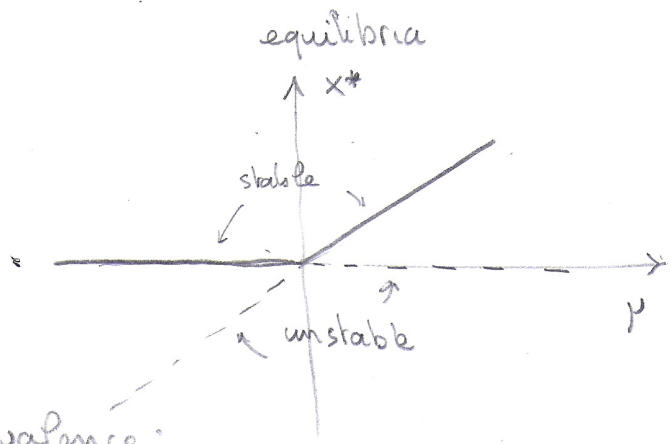
$\dot{X} = F(X, \mu)$ with $F(X_0, \mu_0) = 0$ is locally topologically equivalent to $\dot{x} = \mu - x^2$ if:

- nonhyperbolicity: $D_{X_0} F(X_0, \mu_0)$ has a simple zero eigenvalue with righteigenvector: V
left eigenvector: U
- transversality: $U \cdot \partial_{\mu_0} F(X_0, \mu_0) \neq 0$: the eigenvalue crosses the imaginary axis when μ crosses μ_0
- non degeneracy: $U \cdot D_X^2 F(X_0, \mu_0)(V, V) \neq 0$: the dominant effect is due to quadratic terms

↳ "special" saddle node: different / nondegeneracy transversality conditions
 ↳ transcritical bifurcation: normal form $\dot{x} = \mu x - x^2$



bifurcation diagram

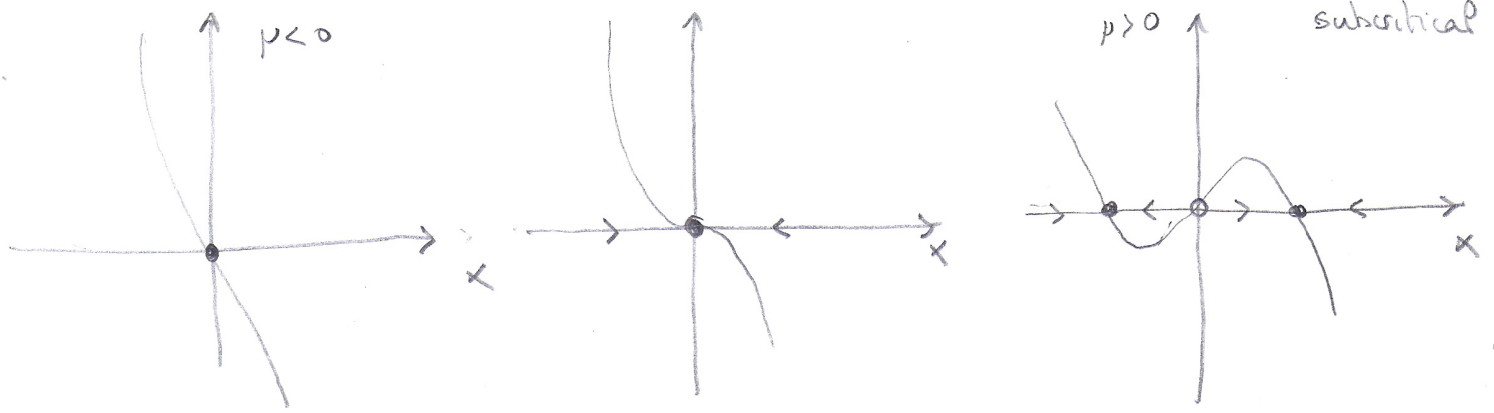


"exchange of" stability

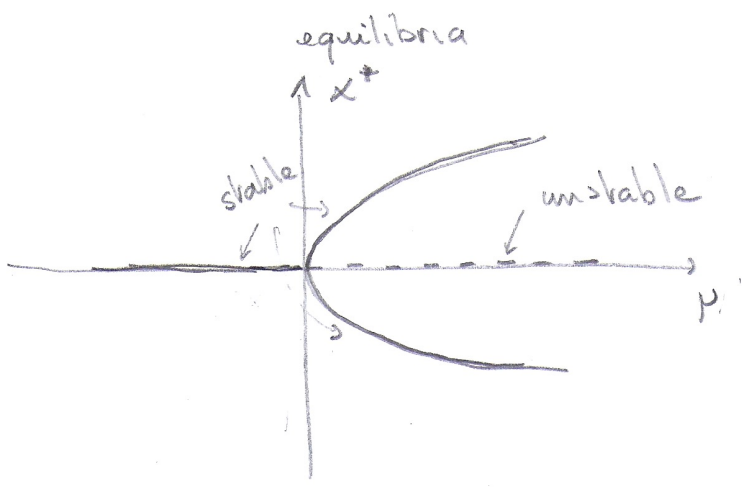
Topological equivalence:

- transversality: under assumption that $\partial_p F(x_0, \mu_0) = 0$ we require $U.D_{\mu x} F(x_0, \mu_0)(V) \neq 0$: x_0 remains an equilibrium.

↳ pitchfork bifurcation, normal form $\dot{x} = \begin{cases} \mu x - x^3 \\ \mu x + x^3 \end{cases}$



"bifurcation"
diagram



Topological equivalence: (typically for $F(x, \mu) = -F(x, \mu)$)

- transversality: under assumption that $\partial_\mu F(x_0, \mu_0) = 0$ we require $U. D_{\mu x} F(x_0, \mu_0)(V) \neq 0$: x_0 remains an equilibrium
- nondegeneracy: under assumption that $U. D_x^2(x_0, \mu_0)(V, V) = 0$, we require $U. D_x^3(x_0, \mu_0)(V, V, V) \neq 0$.

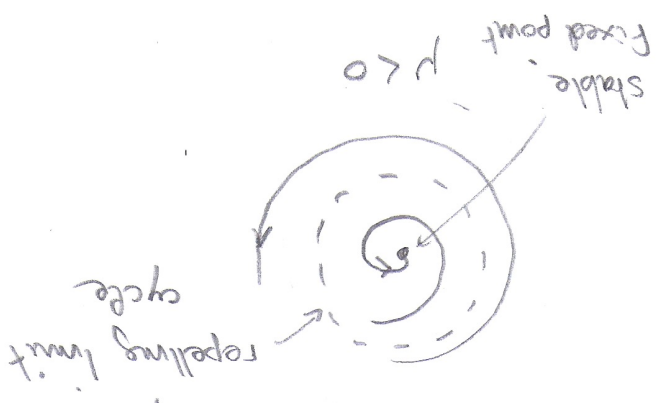
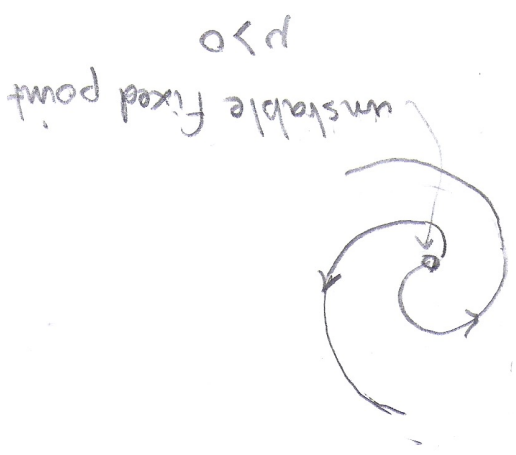
$U. D_x^3(x_0, \mu_0)(V, V, V) < 0 \rightarrow$ supercritical
 $> 0 \rightarrow$ subcritical

* Andronov-Hopf bifurcation: (2-D center manifold)

\hookrightarrow normal form in polar coordinates (infinite number of resonance):

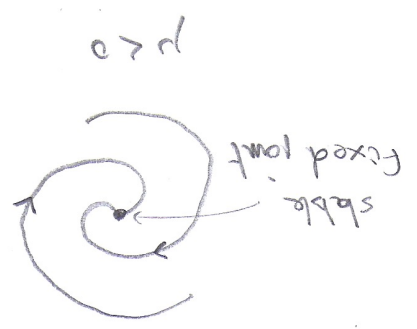
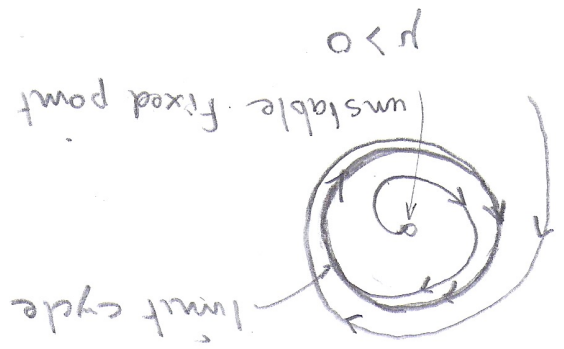
$$\dot{r} = \mu r + a_1 r^3 + a_2 r^5 + a_3 r^7 + \dots$$

$$\dot{\theta} = \omega + b_1 r^2 + b_2 r^4 + b_3 r^6 + \dots$$



$$\begin{cases} \dot{r} = \mu r + r^3 \\ \dot{\theta} = \omega + b r^2 \end{cases}$$
 → "destabilizing" r^3

↳ subcritical Hopf bifurcation $a_1 > 0, a_2 < 0$



$$\begin{cases} \dot{r} = \mu r - r^3 \\ \dot{\theta} = \omega + b r^2 \end{cases}$$
 → "stabilizing" r^3

↳ supercritical Hopf bifurcation $a_1 < 0$

Topological equivalence for Hopf bifurcation

$\dot{X} = F(X, \mu)$ with $F(X_0, \mu_0) = 0$ is locally topologically equivalent to $\dot{r} = (\mu + ar^2)r$, $\dot{\theta} = \omega$ if

- nonhyperbolicity: $D_x F(X_0, \mu_0)$ has a simple pair $(\lambda, \bar{\lambda})$ of pure imaginary eigenvalues and no other eigenvalues have zero real part.
- transversality: $\partial_\mu \operatorname{Re}(\lambda(\mu))|_{\mu=\mu_0} \neq 0$: the pair of eigenvalues cross the imaginary axis when μ crosses μ_0 .
- nondegeneracy: complicated except for 2-D flow it boils down to checking that $a_1 \neq 0$. The dominant effect is due to r^3 .